

**CONDUCTING SYSTEMS OF PLANT LEAVES AS OPTIMAL NETWORKS FOR FLUID DELIVERY**

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Nature inspired solutions for biomimetic design of the technical and biomedical units based on the fluid/gas transportation through a network of ducts/tubes are widely discussed in recent literature [1,2]. The regularities between the lengths  $L$ , diameters  $D$  and bifurcation angles  $\alpha_{1,2}$  of the tubes of different size and branching order (fig.1) can be obtained from the solution of the correspondent optimization problem for a simple cell fed by a single tube of given  $L$  and  $D$  [3].

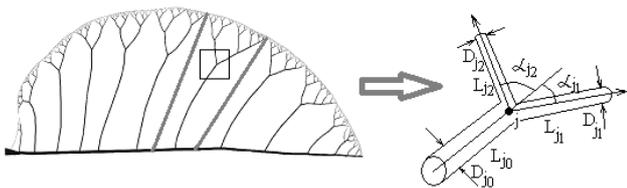


Fig. 1. A sketch of the half leaf blade and the bifurcation of tubes as design unit of the water transportation system

A concept of the microcirculatory cell as a unit of construction of the complex fluid delivery, heat and mass exchange systems with microfluid as a working liquid substance is studied here.

Let us consider a rectangular microcirculatory cell of volume  $V = V_1 + V_2$ , where

$$V_1 = \{x \in [0, L], y \in [-a, a], z \in [-h, h]\},$$

$$V_{II} = \{x \in [0, L], y \in [-H - a, -a] \cup [a, a + H], z \in [-h, h]\},$$

$h \ll \{H, L\}$ ,  $(x, y, z)$  is the Cartesian coordinate system.

The volumes  $V_1$  and  $V_{II}$  are filled with anisotropic porous media with different permeability. The model corresponds to the leaf vein ( $V_1$ ) supplying its drainage area ( $V_{II}$ ) by water and nutrition.

Flow is driven by the hydrostatic pressure difference at the inlet ( $P^+$ ) and outlet ( $P^-$ ) of the tube, as well as the osmotic pressure difference in the sap ( $\pi^+$ ) and in the volume  $V_{II}$  ( $\pi^-$ ). Flow in the cell is governed by Darcy law for viscous incompressible liquid. The governing equations for  $V_1$  and  $V_{II}$  are

$$\text{in } V_1: \quad \text{div } \vec{U} = 0, \quad \vec{U} = -\frac{\hat{K}}{\mu} \nabla P, \quad (1)$$

$$(\vec{U} \nabla) C = -\text{div}(D_c \nabla C) \quad (2)$$

$$\text{in } V_{II}: \quad \text{div } \vec{v} = 0, \quad \vec{v} = -\frac{\hat{M}}{\mu} (\nabla p - \zeta \nabla \pi), \quad (3)$$

$$(\vec{v} \nabla) b = -\text{div}(D_b \nabla b) - q_b \quad (4)$$

where  $\vec{U}$  and  $\vec{v}$ ,  $P$  and  $p$ ,  $\hat{K}$  and  $\hat{M}$ ,  $C$  and  $b$ ,  $D_c$  and  $D_b$  are the flow velocities, hydrostatic pressures, permeability tensors for the liquid, concentrations and

diffusion coefficients in  $V_1$  and  $V_{II}$  accordingly;  $\pi$  is the osmotic pressure in  $V_{II}$ ;  $\mu$  is the fluid viscosity;  $q_b$  is the absorption rate of the dissolved substances in the volume (catalytic layer, live cells, etc);  $\zeta$  is the selectivity coefficient governing the relative conductivity for the dissolved substances (provided they are large), so  $\zeta \hat{M}$  is conductivity for the chemical components.

The boundary conditions with the Maxwell first order velocity slip model at the wall within the Knudsen layer are assumed in the form [4]

$$\begin{aligned} \text{in } V_1: \quad x=0: \quad C=C^+, \quad P=P^+, \\ x=L: \quad C=C^-, \quad P=P^-, \\ z=0; \quad y=0: \quad U_y=0, \quad U_z=0, \end{aligned} \quad (5)$$

$$y=\pm a: \quad U_{x,z} = \frac{2-\sigma_{x,z}}{\sigma_{x,z}} \text{Kn} \frac{\partial U_{x,z}}{\partial U_{x,z}}, \quad U_y = \pm V_f,$$

$$z=\pm h: \quad U_z = w_{\pm}^I.$$

$$\begin{aligned} \text{in } V_{II}: \quad x=0;L: \quad \frac{\partial b}{\partial x} = 0, \quad \frac{\partial p}{\partial x} = 0, \\ y=\pm(a+H): \quad \frac{\partial b}{\partial y} = 0, \quad \frac{\partial p}{\partial y} = 0, \end{aligned} \quad (6)$$

$$z=\pm h: \quad v_z = w_{\pm}^{II}, \quad \frac{\partial b}{\partial z} = \alpha b w_{\pm}^{II},$$

where  $C^{\pm}$  are the concentrations at the inlet and outlet of the duct  $V_1$ ;  $w_{\pm}^I$  and  $w_{\pm}^{II}$  are the mass exchange rates per unit surface area at the upper (+) and lower (-) surfaces of  $V_1$  and  $V_{II}$ ;  $\alpha = \text{const}$  corresponds to washing out the unabsorbed components,  $\alpha = 0$  when the outer walls of the cell are impermeable for the components,  $n$  is the normal direction (coordinate) to the wall,  $\text{Kn} = \lambda / L$  is the Knudsen number,  $\lambda$  is the mean free path of the microparticle in the fluid,  $L$  is the length,  $\sigma$  is the tangential momentum accommodation factor (or friction);  $\sigma = 1$  for purely diffuse reflection.

The lumped parameter model have been obtained from (1)-(6) by averaging over  $y, z$ . The solution has been found in analytical form and analyzed. The problem of hydrodynamically interacting cells is studied. Possible applications in the technical heat and mass exchangers, microheaters/coolers and fuel cells are discussed.

LITERATURE

1. Kizilova N., Szekeres A. Nature Inspired Optimal Composites. In: Continuum physics and engineering applications. – 2010. – P. 30–42.
2. Hamadiche M., Kizilova N. Nature inspired optimal design of heat conveying networks for advanced fiber-reinforced composites. // J. Thermal Eng. – 2015. – V.1, Issue 7. – P. 636–645.
3. Kizilova N. Liquid filtration in a microcirculatory cell of the plant leaf. // Intern. J. Fluid Mech. Res. – 2007. – v.34, N6. – P.572–588.
4. Karniadakis G.E., Beskok A., Aluru N. Microflows and nanoflows: Fundamentals and simulation. Interdisc. Appl. Math. Series, V.29. 2005. P.51–77.